Langevin description for evaporation-fission processes in CRISP code

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Introduction

2 Langevin equation

- Parametrization
- Langevin equation
- Entropy and temperature
- Deformation potential
- Emission width

3 Results and discussion



Langevin dynamic:

- The time evolution of the nucleus is considered and friction β parameter that has a direct influence in the fission probability and evaporation multiplicity.
- This study is more interesting for ion reaction, where a considerable angular momentum is transfered to compound nucleus.

• The following simple polynomial expression is in rather good agreement with the known results of exact LDM calculation:

$$\pi(u, v) = v^2 - (1 - u)^2 (A + Bu^2 + \alpha u)$$
(1)

• Volume conservation:

$$c = (A + \frac{1}{5}B)^{-1/3}$$
 (2)

• Connection between the parameter sets A,B and c,h:

$$B = 2h + 1/2(c - 1)$$
(3)

$$A = \frac{1}{c^3} - \frac{1}{5}B \tag{4}$$

Parametrization (c,h)



• Langevin equation for over-damped motion:

$$\dot{q} = -rac{1}{Meta}rac{dV(q)}{dq} + \sqrt{rac{T}{Meta}}\Gamma(t)$$
 (5)

• Langevin discretized form:

$$q_{n+1} = q_n + \left[\sqrt{\frac{T(q)}{\beta(q)M}} \frac{dS(q)}{dq}\right]_n \tau + \sqrt{\left[\frac{T(q)}{\beta(q)M}\right]_n \tau} \omega_n \qquad (6)$$
$$q(c,h) = \frac{3}{8}c(1 + \frac{2}{15}c^3)$$

• Fermi's gas expression for the entropy:

$$S(q, E_{tot}^*, A, Z, L) = 2\sqrt{a(q, A)[E_{tot}^* - V(q, A, Z, L)]}$$
(7)

• The temperature is:

$$T(q, E_{tot}^*, A, Z, L) = \frac{S(q, E_{tot}^*, A, Z, L)}{2a(q, A)}$$
(8)

• The potential energy is given by the liquid-drop model:

$$V_{s}(q, A, Z) = a_{2} \left[1 - k \left(\frac{N - Z}{A} \right)^{2} \right] A^{2/3} \left[B_{s}(q) - 1 \right]$$
(9)

$$B_s(q) = \begin{cases} 1 + 0.4(64/9)(q - 0.375)^2 & \text{if } q < 0.452 \\ 0.983 + 0.439(q - 0.375) & \text{if } q \ge 0.452 \end{cases}$$
(10)

$$a_2 = 17.9439 MeV \quad k = 1.7826$$
(11)

Deformation potential

• The potential energy is given by the liquid-drop model:

$$V_c(q, A, Z) = c_3 \frac{Z^2}{A^{1/3}} [B_s(q) - 1]$$
(12)

$$B_{c}(x) = \begin{cases} 1 + \left(1 - B_{s} + \frac{B_{f}}{E_{ssp}}\right)/2X & \text{if } q < 0.452\\ 1 - 0.2\frac{64}{9}(q - 0.375)^{2} & \text{if } q \ge 0.452 \end{cases}$$
(13)

$$\frac{B_f}{E_{ssp}}(q) = \begin{cases} 0.2599 - 0.2151X - 0.1643X^2 \\ -0.0673X^3 & \text{if } q < 0.6 \\ 0.7259Y^3 - 0.3302Y^4 + 0.6387Y^5 \\ +7.827Y^6 - 12.0061Y^7 & \text{if } q \ge 0.452 \end{cases}$$

$$X = 0.05 \ln[0.875(q_{sd} - 0.375)^{-1} - 1] + 0.74$$
(14)

Image: A matrix

Deformation potential

• The potential energy is given by the liquid-drop model:

$$B_r = J_{||}^{-1}$$
 if $J_\perp < J_{||}$ and $q > 0.375$
 $B_r = J_\perp^{-1}$ in all other cases

$$J_{\parallel} = c^{2} \left\{ c^{-3} + 4B_{sh} [(4/15)B_{sh}c^{3} - 1]/35 \right\}$$
$$J_{\perp} = c^{2} \left\{ 1 - c^{-3} + 4B_{sh} [(4/15)B_{sh}c^{3} - 1]/35 \right\} / 2$$

$$B_{sh}(c,h) = 2h + (c-1)/2$$
$$q(c,h) = \frac{3}{8}c(1 + \frac{2}{15}c^3)$$

Image: A matrix



• Woods-Saxon single particle potential:

$$a(q, A) = \alpha_1 A + \alpha_2 A^{2/3} B_s(q)$$

 $\alpha_1 = 0.073 \ MeV^{-1} \quad \alpha_2 = 0.095 \ MeV^{-1}$

• The Standard Parameter Set friction

$$\beta_{SPS}(q) = \begin{cases} \beta_0 & \text{if } q < q_{neck} \\ \beta_0 + \frac{\beta_{sc} - \beta_0}{q_{sc} - q_{neck}} (q - q_{neck}) & \text{if } q_{neck} \le q \le q_{sc} \end{cases}$$



• The emission width of a particle of kind ν ($\nu = n, p, \alpha$):

$$\Gamma_{\nu} = (2s_{\nu}+1)\frac{m_{\nu}}{\pi^{2}\hbar^{2}\rho_{c}(E^{*})}\int_{0}^{E^{*}-B_{\nu}}d\varepsilon_{\nu}\rho_{R}(E^{*}-\varepsilon_{\nu})\varepsilon_{\nu}\sigma_{inv}(\varepsilon_{\nu})$$

• Level densities of the compound and residual nuclei are:

$$ho \sim (2L+1) exp(S)$$

Inverse cross sections are:

$$\sigma_{in\nu}(\varepsilon_{\nu}) = \begin{cases} \pi R_{\nu}^{2}(1 - V_{\nu}/\varepsilon_{\nu}) & \text{for } \varepsilon_{\nu} > V_{\nu} \\ 0 & \text{for } \varepsilon_{\nu} < V_{\nu} \end{cases}$$

• The barriers for the charged particles are:

$$V_{
u} = rac{(Z - Z_{
u})Z_{
u}K_{
u}}{R_{
u} + 1.6}$$
 $R_{
u} = 1.22 \left[(A - A_{
u})^{1/3} + A_{
u}^{1/3}
ight] + rac{3.4}{arepsilon_{
u}^{1/2} \delta_{
u,n}}$

 ${\it K}_{\nu}=1.32$ for lpha ${\it K}_{
u}=1.15$ for proton

• For giant dipole $\gamma - quanta$:

$$\Gamma_{\gamma} = \frac{3}{\rho_{C}(E^{*})} \int_{0}^{E^{*}} d\varepsilon \rho_{C}(E^{*} - \varepsilon) f(\varepsilon)$$
$$f(\varepsilon) = \frac{4}{3\pi} \frac{1 + k}{m_{n}c^{2}} \frac{e^{2}}{\hbar c} \frac{NZ}{A} \frac{\Gamma_{G}\varepsilon^{4}}{(\Gamma\varepsilon)^{2} + (\varepsilon^{2} - E_{G}^{2})^{2}}$$

with k=0.75, E_G =80/ $A^{1/3}$ MeV, Γ_G =5 MeV

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• For fission:

$$\Gamma_{f} = \hbar \frac{T_{gs}}{\beta_{gs}} \frac{\sqrt{\left|S_{gs}^{''}\right|S_{sd}^{''}}}{\pi M} \exp(S(q_{sd}) - S(q_{gs})) \times \left\{1 + erf\left[(q_{sc} - q_{sd})\sqrt{S_{sd}^{''}/2}\right]\right\}^{-1}$$

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• The total initial excitation energy is:

$$E_{tot}^* = E_{cm} + Q$$

where:

$$E_{cm} = E_{lab}A_T / (A_T + A_P)$$

• The angular momentum $L = \hbar I$ is sampled from the spin distribution:

$$\sigma(l) = \frac{2\pi}{k^2} \frac{2l+1}{1 + \exp\left[(l - l_c)/\delta l\right]}$$

• The quantity I_c scales as

$$I_{c} = \sqrt{A_{p}A_{T}/A} \left(A_{p}^{1/3} + A_{T}^{1/3} \right) \left(0.33 + 0.205 \sqrt{E_{cm} - V_{c}} \right)$$

• The barrier V_c is:

$$V_{c} = 5/3c_{3}Z_{P}Z_{T} / \left(A_{P}^{1/3} + A_{T}^{1/3} + 1.6\right)$$

with $c_3 = 0.7053 \text{ MeV}$

• The diffuseness δI is:

$$\left\{ \begin{array}{ll} (A_P A_T)^{3/2} * 10^{-5} \left[1.5 - 0.02 \left(E_{cm} - V_c - 10 \right) \right] & \text{ for } E_{cm} > V_c + 10 \\ (A_P A_T)^{3/2} * 10^{-5} \left[1.5 - 0.04 \left(E_{cm} - V_c - 10 \right) \right] & \text{ for } E_{cm} < V_c + 10 \end{array} \right.$$



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Conclusions

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